

A Least-Squares Finite Element formulation of the Geometrically-Nonlinear Elasticity Equations

Chad Westphal

`westphac@wabash.edu`

Department of Math and CS

Wabash College

Crawfordsville, IN

Outline

- Introduction
 - First-Order System Least Squares (FOSLS) methodology
- Elasticity
 - Elasticity notation, definitions and FOSLS formulation
 - Theory overview
 - Computational results: Smooth solutions
 - Regularity, what (not) to expect
- FOSLS in Nonsmooth Domains
 - Weighted spaces and the weighted-norm approach
 - Theory and error bounds
 - Computational results: Nonsmooth solutions
- Concluding remarks and questions

Basic Least-Squares Methodology

Consider solving the first-order system $L\mathbf{u} = \mathbf{f}$ by minimizing

$$G(\mathbf{u}; f) = \|L\mathbf{u} - \mathbf{f}\|^2$$

over all $\mathbf{u} \in \mathcal{V}$. This leads to the weak problem: find $\mathbf{u} \in \mathcal{V}$ such that

$$(L\mathbf{u}, L\mathbf{v}) = (\mathbf{f}, L\mathbf{v}) \quad \forall \mathbf{v} \in \mathcal{V}.$$

Basic Least-Squares Methodology

Consider solving the first-order system $L\mathbf{u} = \mathbf{f}$ by minimizing

$$G(\mathbf{u}; f) = \|L\mathbf{u} - \mathbf{f}\|^2$$

over all $\mathbf{u} \in \mathcal{V}$. This leads to the weak problem: find $\mathbf{u} \in \mathcal{V}$ such that

$$(L\mathbf{u}, L\mathbf{v}) = (\mathbf{f}, L\mathbf{v}) \quad \forall \mathbf{v} \in \mathcal{V}.$$

\mathcal{V} ellipticity: If $G(\mathbf{u}; 0)^{\frac{1}{2}} = \|L\mathbf{u}\|$ is equivalent to the norm on \mathcal{V} :

$$c_0 \|\mathbf{u}\|_{\mathcal{V}} \leq \|L\mathbf{u}\| \leq c_1 \|\mathbf{u}\|_{\mathcal{V}},$$

then the weak problem is well-posed in \mathcal{V} .

Discrete Problem

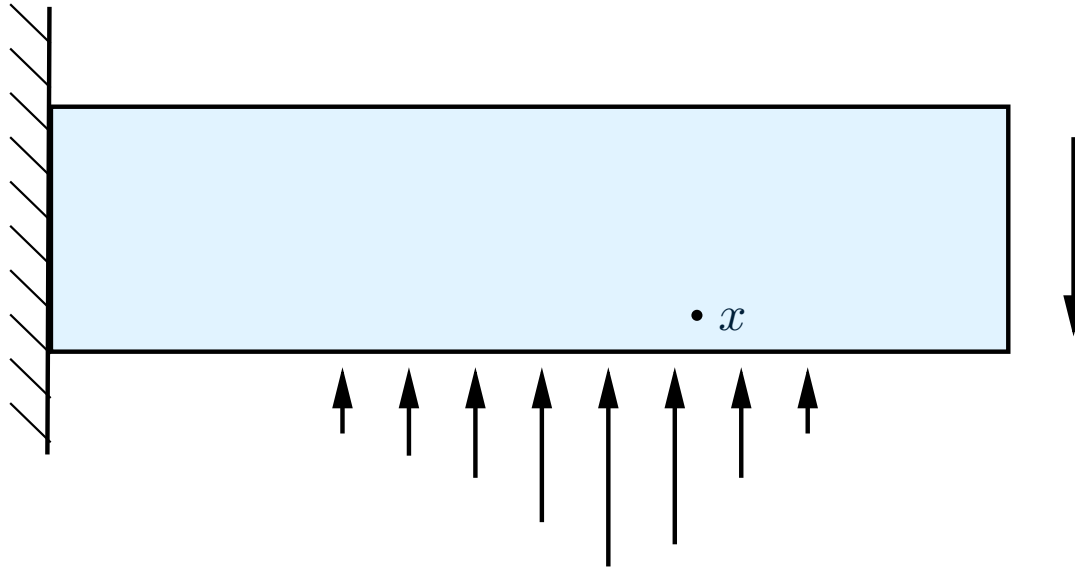
Consider the FE space $\mathcal{V}^h \subseteq \mathcal{V}$. The discrete weak problem is: find $\mathbf{u}^h \in \mathcal{V}^h$ such that

$$(L\mathbf{u}^h, L\mathbf{v}^h) = (\mathbf{f}, L\mathbf{v}^h) \quad \forall \mathbf{v}^h \in \mathcal{V}^h.$$

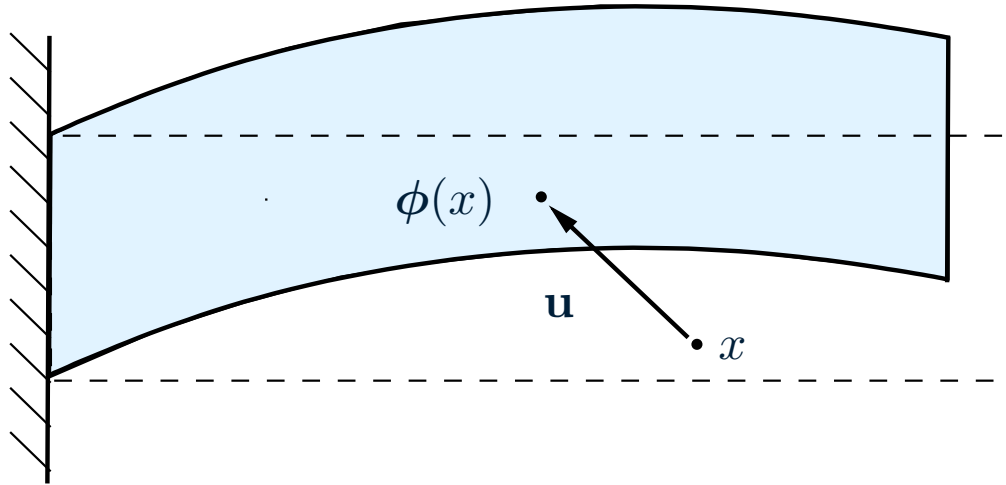
When $\mathcal{V}^h \subseteq H^1(\Omega)$ and $\|\cdot\|_{\mathcal{V}} = \|\cdot\|_1$, H^1 elliptic functionals result in:

- Existence and uniqueness of $\mathbf{u}^h \in \mathcal{V}^h$
- Optimal order FE discretization accuracy (given regularity of $L\mathbf{u} = \mathbf{f}$)
- $O(h^{-2})$ conditioned linear systems solved with standard multigrid methods
- Local functional values a free H^1 *a posteriori* error bound

Elasticity Notation and Definitions



Elasticity Notation and Definitions



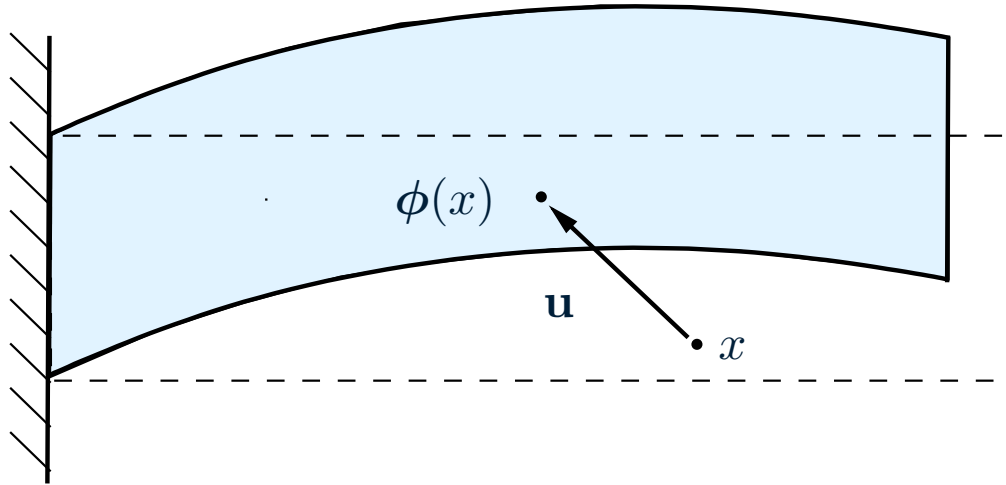
\mathbf{u}	Displacement
ϕ	Deformation
E	Strain Tensor
Σ	Stress Tensor

$$\phi(x) = x + \mathbf{u}(x)$$

$$\begin{aligned} E &= \frac{1}{2}(I - \nabla\phi^t\nabla\phi) \\ &= \frac{1}{2}(\nabla\mathbf{u} + \nabla\mathbf{u}^t + \nabla\mathbf{u}^t\nabla\mathbf{u}) \end{aligned}$$

$$\Sigma = \lambda \text{tr}(E)I + 2\mu E + o(E)$$

Elasticity Notation and Definitions



\mathbf{u}	Displacement
ϕ	Deformation
E	Strain Tensor
Σ	Stress Tensor

$$\phi(x) = x + \mathbf{u}(x)$$

$$\begin{aligned} E &= \frac{1}{2}(I - \nabla\phi^t\nabla\phi) \\ &= \frac{1}{2}(\nabla\mathbf{u} + \nabla\mathbf{u}^t + \nabla\mathbf{u}^t\nabla\mathbf{u}) \end{aligned}$$

$$\Sigma = \lambda \text{tr}(E)I + 2\mu E + o(E)$$

Nonlinear Elasticity Equations

$$\left\{ \begin{array}{ll} \nabla \cdot [(I + \nabla \mathbf{u})\Sigma(\nabla \mathbf{u})] = \mathbf{f} & \text{in } \Omega \\ \mathbf{n} \cdot (I + \nabla \mathbf{u})\Sigma(\nabla \mathbf{u}) = \mathbf{g} & \text{on } \Gamma_T \\ \mathbf{u} = \mathbf{0} & \text{on } \Gamma_D \end{array} \right.$$

Nonlinear Elasticity Equations

$$\left\{ \begin{array}{ll} \nabla \cdot [(I + \nabla \mathbf{u})\Sigma(\nabla \mathbf{u})] = \mathbf{f} & \text{in } \Omega \\ \mathbf{n} \cdot (I + \nabla \mathbf{u})\Sigma(\nabla \mathbf{u}) = \mathbf{g} & \text{on } \Gamma_T \\ \mathbf{u} = \mathbf{0} & \text{on } \Gamma_D \end{array} \right.$$

The equilibrium equation may be written as

$$\underbrace{\mu \Delta \mathbf{u} + (\lambda + \mu) \nabla \nabla \cdot \mathbf{u}}_{\text{linear}} + \underbrace{\nabla \cdot P_3(\nabla \mathbf{u})}_{\text{nonlinear}} = \mathbf{f}$$

- Accurate for deformations of “large-displacement and small-strain”
- May scale with $\mu = 1$, and let λ determine incompressibility

Newton-FOSLS

- Pose in terms of $\mathbf{U} = \nabla \mathbf{u}$ (First-Order System)
 - $\Phi = \mathbf{I} + \mathbf{U}$ is the Jacobian of the mapping
 - $\det(\Phi) \approx 1, \|\Phi\|_{\ell^2}^2 \approx 2$
- 2-Stage Solution Algorithm
 - Solve for \mathbf{U}
 - Recover \mathbf{u} from: $\nabla \mathbf{u} = \mathbf{U}$
- Linearize around a current approximation by Newton's method
 - Initial guess: $\mathbf{U}_0 = 0$... the reference configuration
 - General Newton Step: $\nabla \cdot A_n \mathbf{U} = \mathcal{F}_n$
 - 1st step is linear elasticity

Problem Modification

- **Problem:** System matrix $A_n = A(\Phi_n)$ is generally indefinite
- Can modify A_n without changing the solution

Choose B to be such that $\nabla \cdot B = \nabla \times,$

$$\implies \nabla \cdot B \nabla p = 0 \text{ for any } p,$$

Choose $\tilde{A}_n = A_n + B$ then,

$$\nabla \cdot \tilde{A}_n \mathbf{U} = \nabla \cdot A_n \mathbf{U}$$

- Choose B to shift the spectrum of A_n positive...

Properties of \tilde{A}_n

- The geometrically-nonlinear model assumes $o(E) \approx 0$. Thus we may also impose an assumption on the size of $E = \frac{1}{2}(\Phi^t \Phi - I)$.
- The ℓ^2 norm of the strains and the full spectrum of \tilde{A}_n can be written in terms of $\delta = \det(\Phi)$, $\sigma = \|\Phi\|_{\ell^2}^2$.
- We show \tilde{A}_n positive definite under the small strains assumption:

$$\|\Phi^t \Phi - I\|_{\ell^2} < \frac{\sqrt{2}}{\lambda + 3} \implies \tilde{A}_n \text{ is positive definite.}$$

Newton-FOSLS

- Define the FOSLS functional for each step:

$$G(\mathbf{U}; \mathbf{U}_n, \mathcal{F}_n) = \|\nabla \cdot \tilde{A}_n \mathbf{U} - \mathcal{F}_n\|_0^2 + \|\nabla \times \mathbf{U}\|_0^2,$$

for \mathbf{U} in the space

$$\mathcal{V} = \{\mathbf{V} \in H^1(\Omega)^4 : \boldsymbol{\tau} \cdot \mathbf{V} = \mathbf{0} \text{ on } \Gamma_D\}.$$

- Solve by minimization principle: Choose $\mathbf{U} \in \mathcal{V}$ so that

$$G(\mathbf{U}; \mathbf{U}_n, \mathcal{F}_n) = \inf_{\mathbf{V} \in \mathcal{V}} G(\mathbf{V}; \mathbf{U}_n, \mathcal{F}_n).$$

- We show a well-posed Newton minimization, converging by an inexact nested iteration-Newton-FOSLS-multigrid process (cf. Codd, Manteuffel and McCormick).

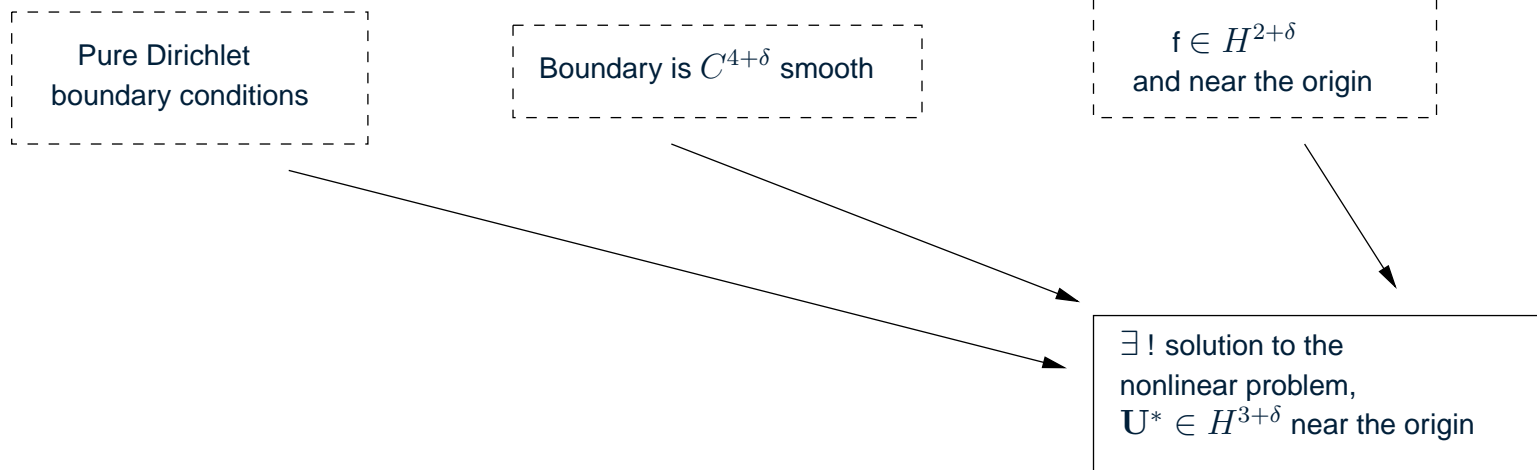
Elasticity Theory

Pure Dirichlet
boundary conditions

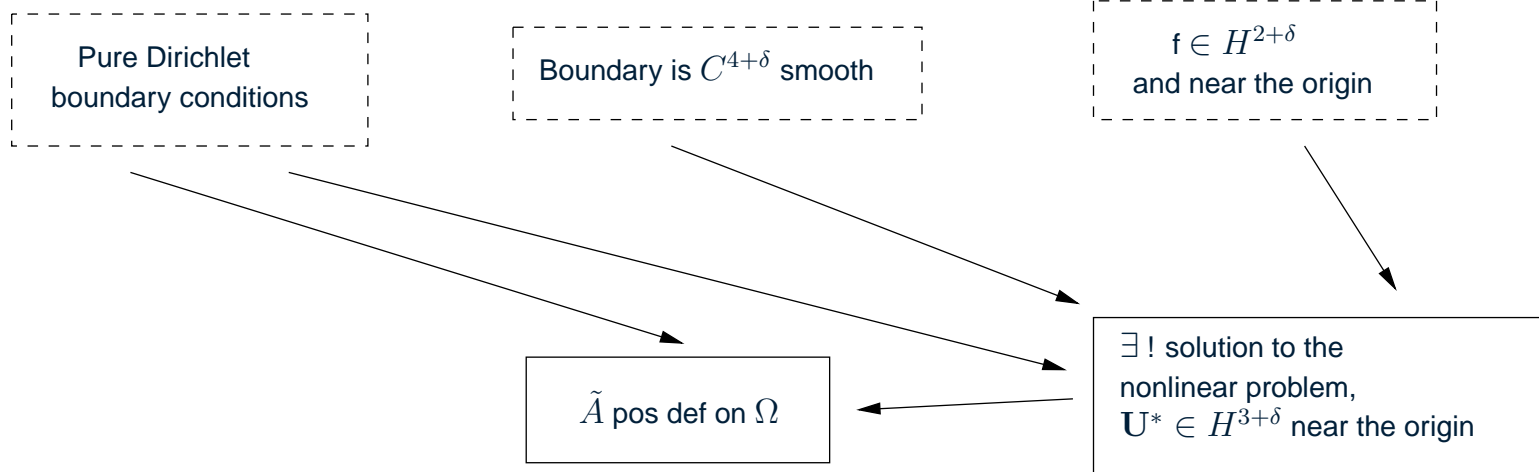
Boundary is $C^{4+\delta}$ smooth

$f \in H^{2+\delta}$
and near the origin

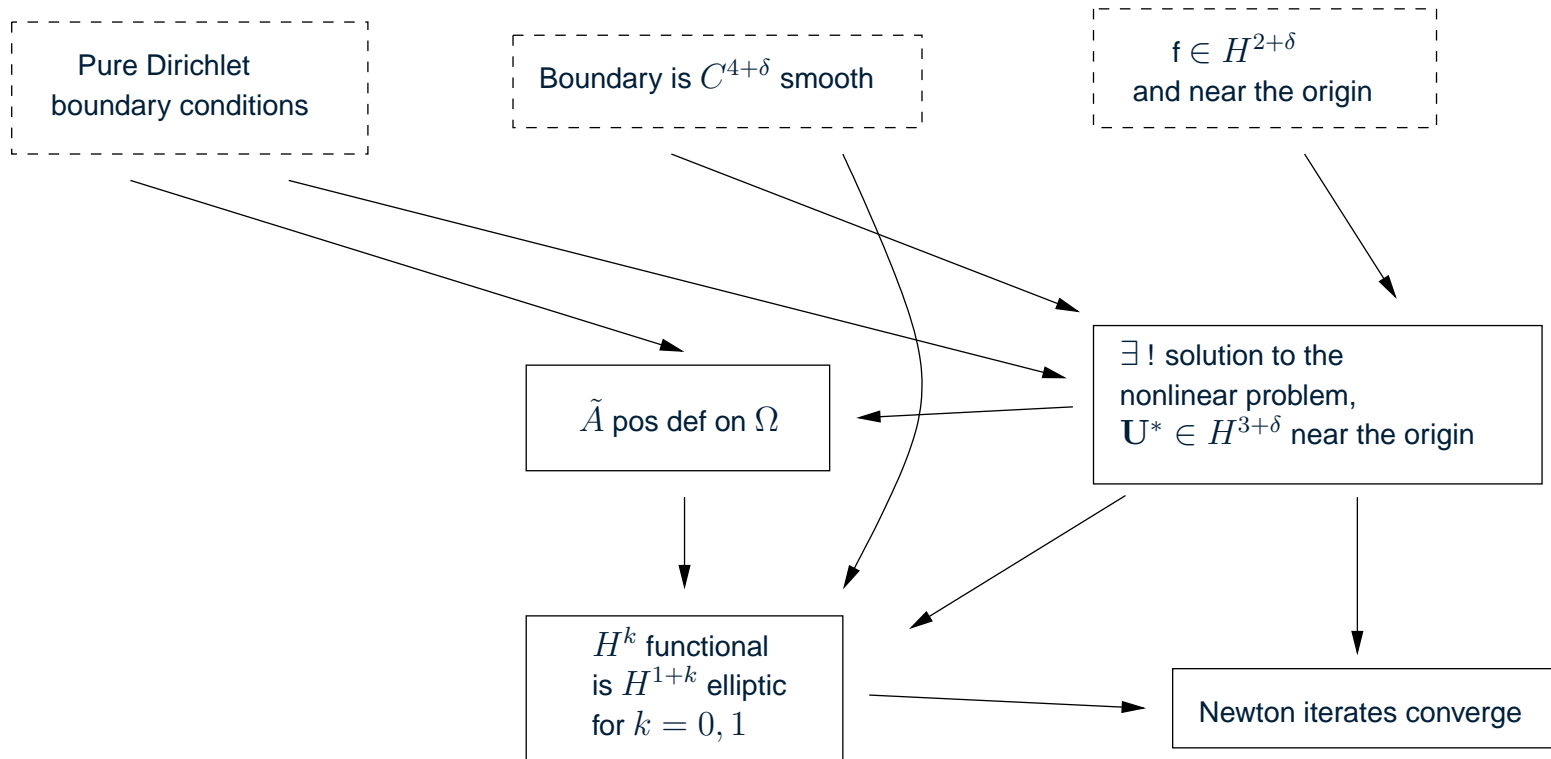
Elasticity Theory



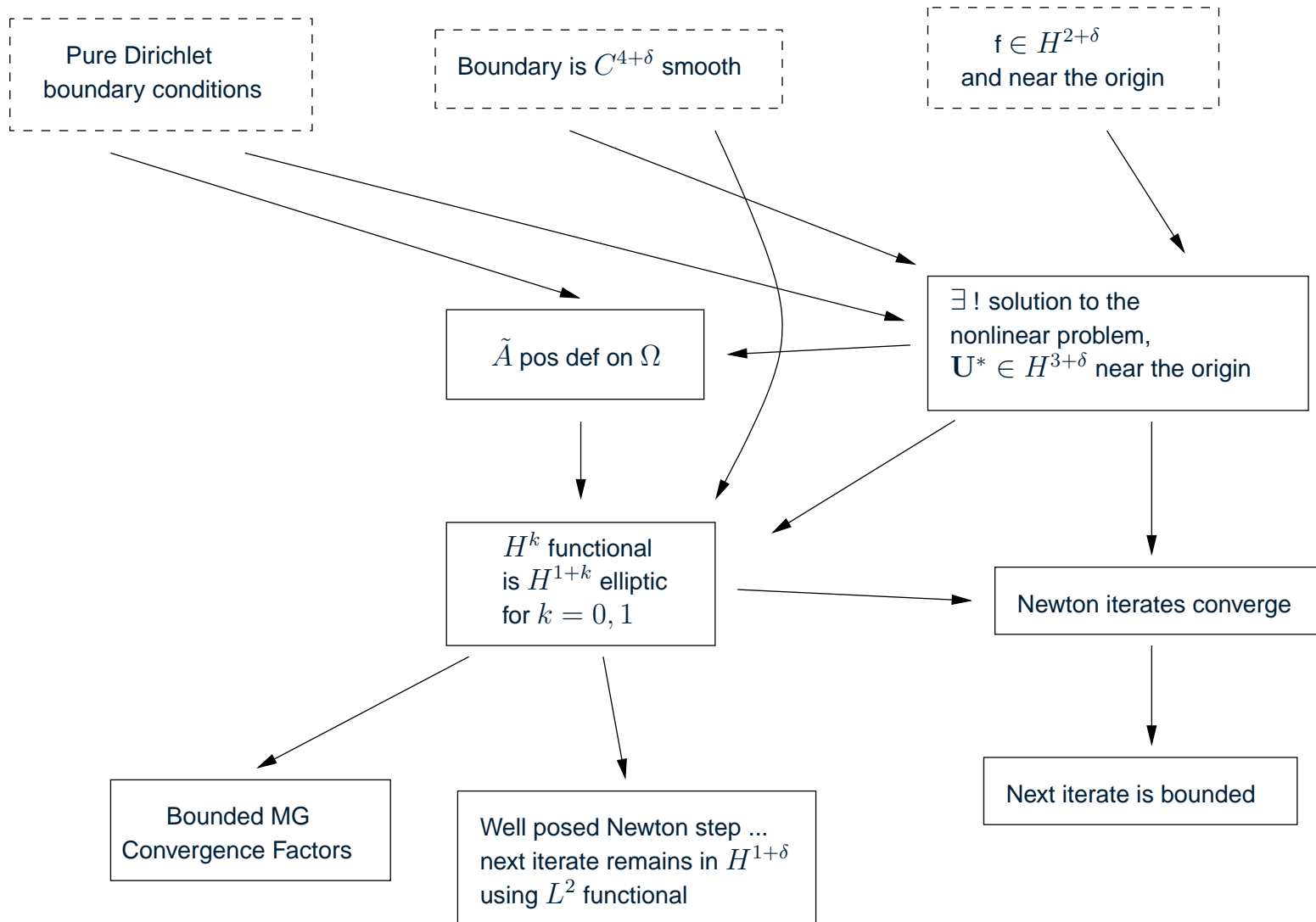
Elasticity Theory



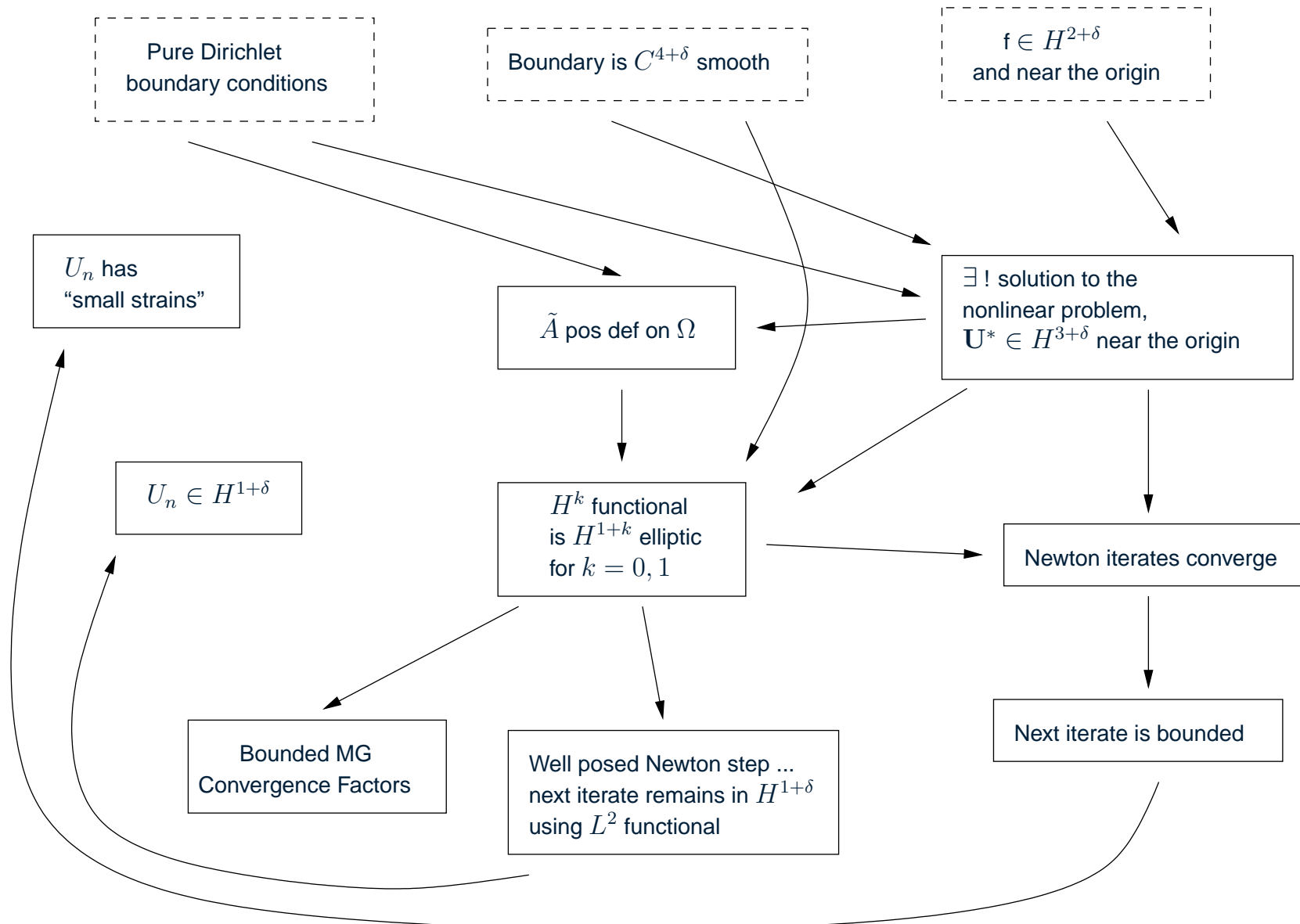
Elasticity Theory



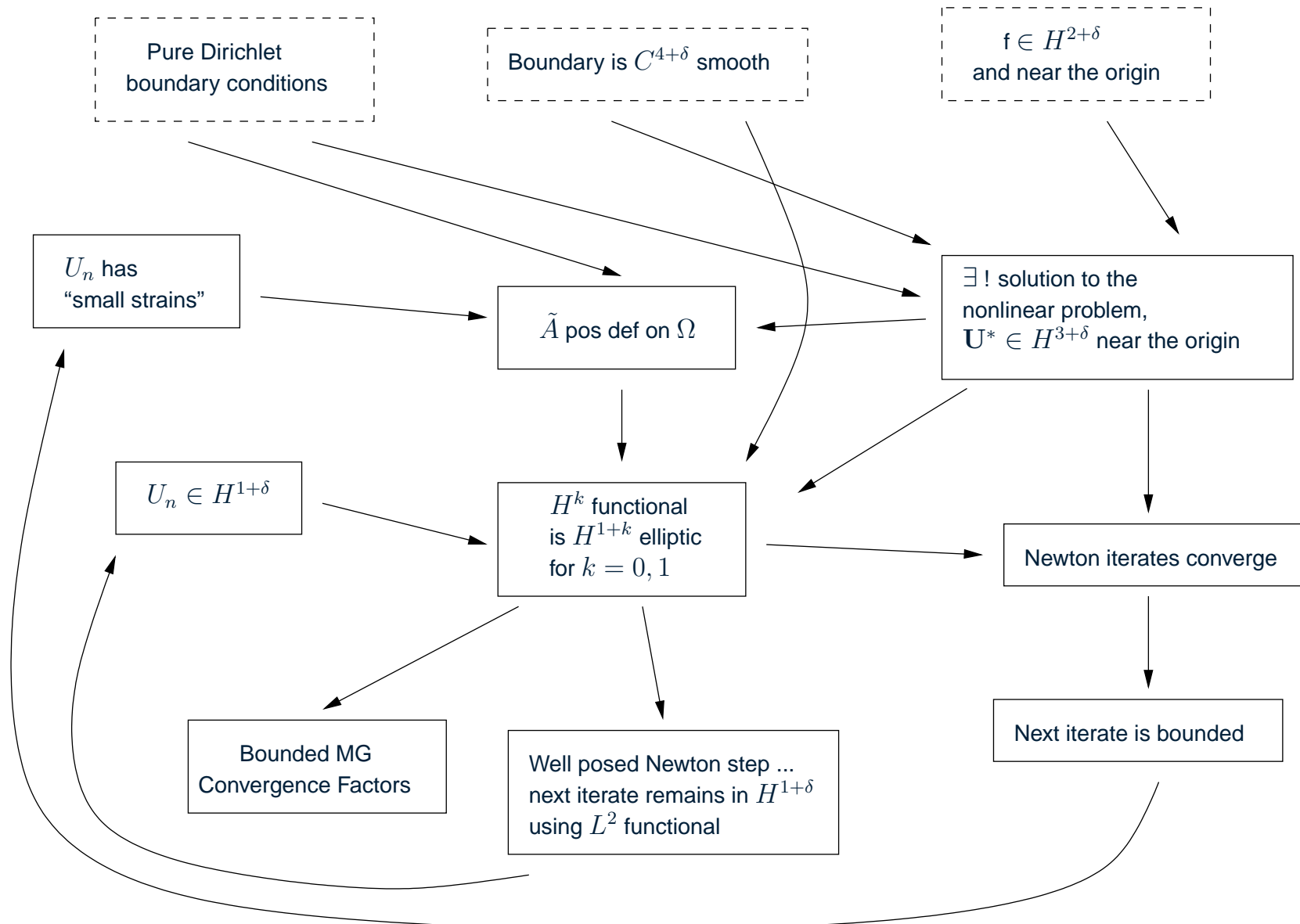
Elasticity Theory



Elasticity Theory



Elasticity Theory



Main Assumptions

- Pure displacement boundary conditions...
 - Mixed and pure traction cases perform similarly in practice.
 - Most likely a sufficient but not necessary condition.

Main Assumptions

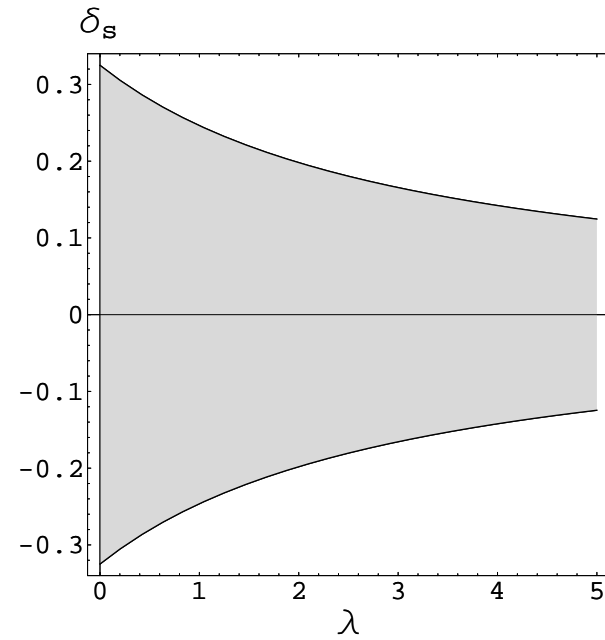
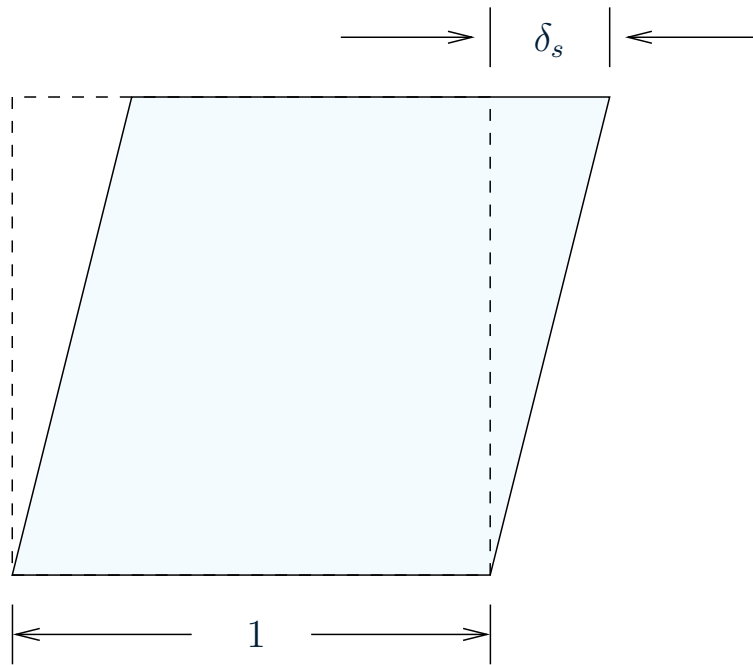
- Pure displacement boundary conditions...
 - Mixed and pure traction cases perform similarly in practice.
 - Most likely a sufficient but not necessary condition.
- The solution has small strains...
 - We seek to show that "small strains" reasonably corresponds to the limits of the model.
 - Most likely a sufficient but not necessary condition.

Main Assumptions

- Pure displacement boundary conditions...
 - Mixed and pure traction cases perform similarly in practice.
 - Most likely a sufficient but not necessary condition.
- The solution has small strains...
 - We seek to show that "small strains" reasonably corresponds to the limits of the model.
 - Most likely a sufficient but not necessary condition.
- The domain is sufficiently smooth...
 - We seek a method for LS problems in nonsmooth domains.
 - A necessary and sufficient condition for most problems.

Small Strains: Pure Shear and Pure Tensile Strain

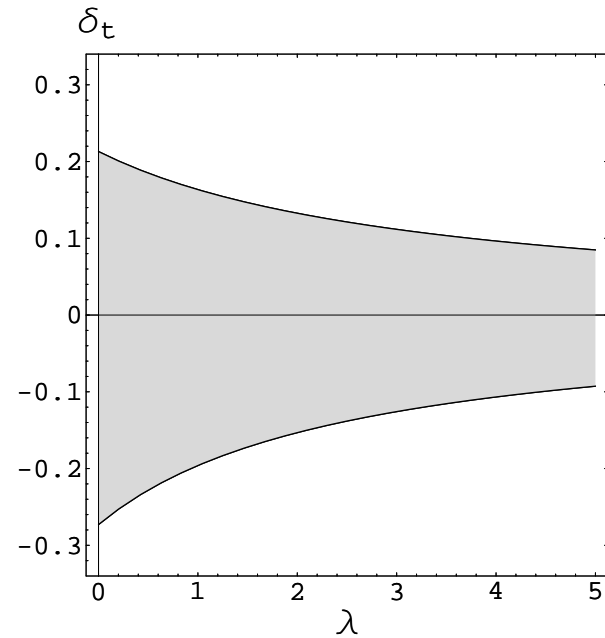
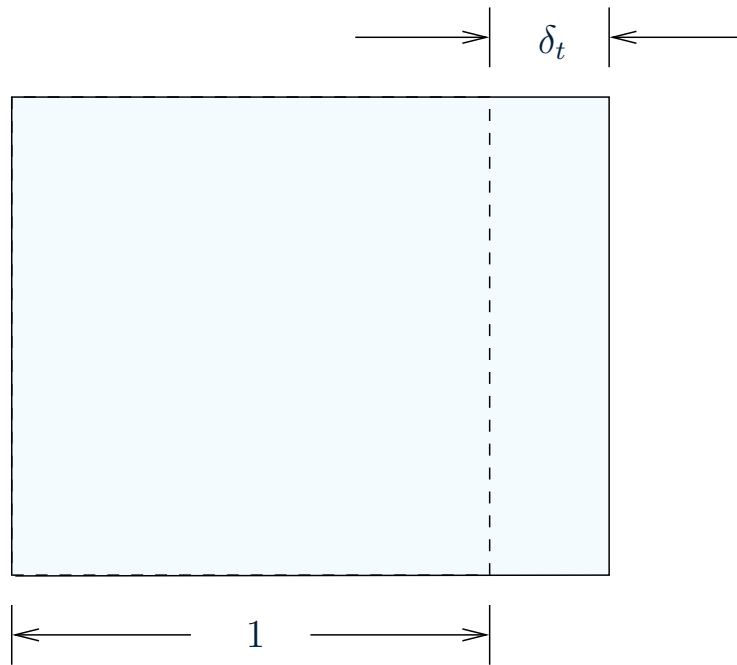
Pure shear strain:



"Small strain" limit of pure shear strain versus λ

Small Strains: Pure Shear and Pure Tensile Strain

Pure tensile strain:



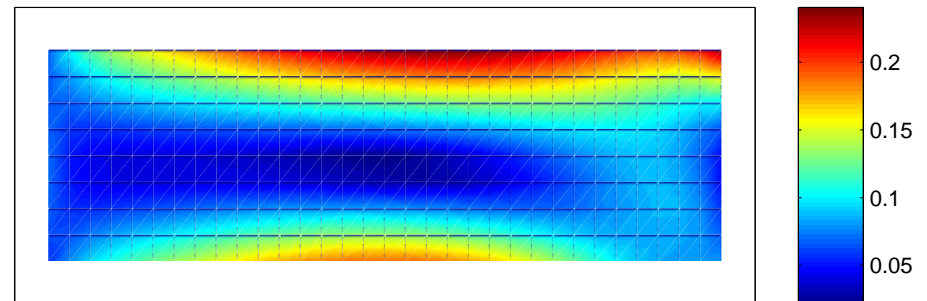
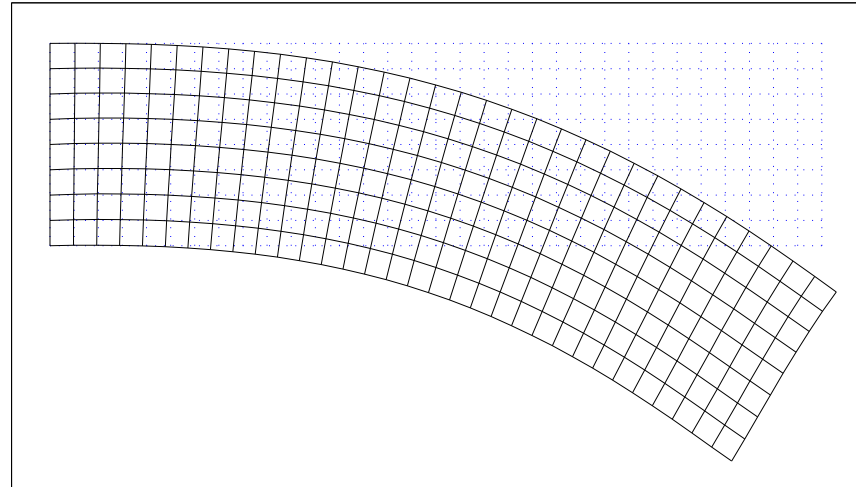
"Small strain" limit of pure tensile strain versus λ

Small strains example

For this configuration our theory holds for $\lambda < 3.25$ ($\nu < 0.38$).

$$\max_{\Omega} \|\Phi^t \Phi - I\|_{\ell^2} = 0.241$$

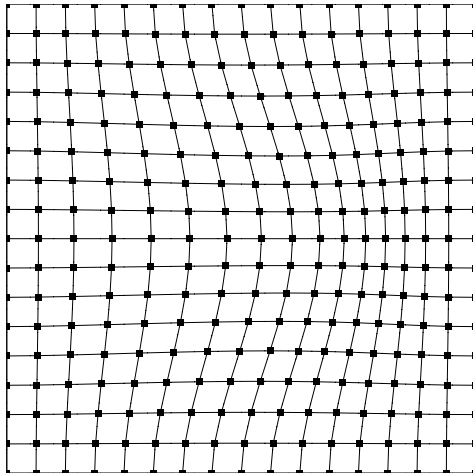
	λ	* ν
Steel:	1.22	0.28
Nickel:	1.53	0.30
Aluminum:	2.15	0.34
maximum:	3.25	0.38
Lead:	7.30	0.44
Rubber:	33.3	0.49



*The Poisson ratio, $\nu = \frac{\lambda}{2(\lambda+1)}$, is another measure of incompressibility.

Numerical Results: Smooth Solutions

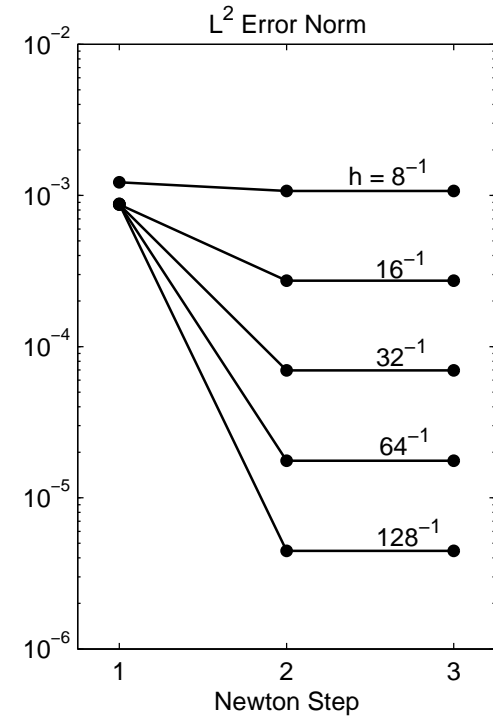
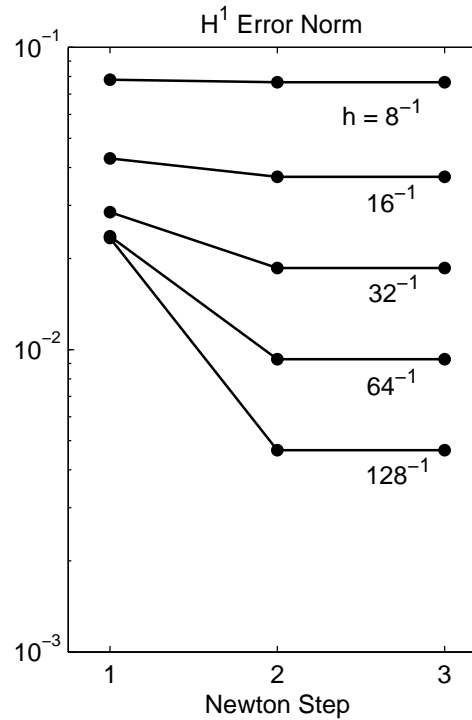
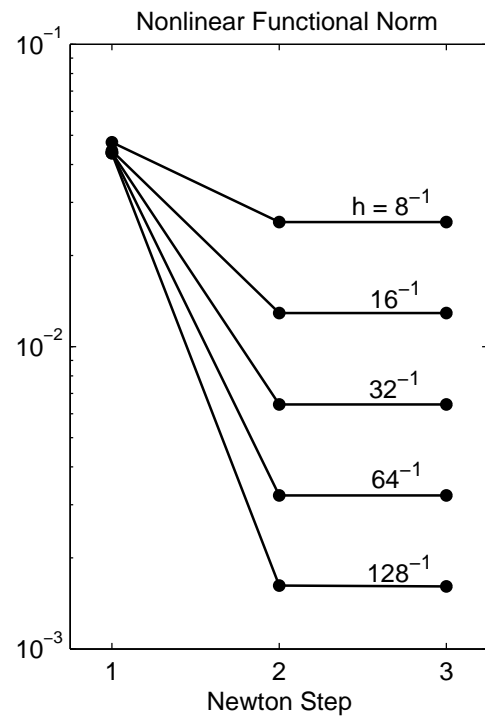
- Displacement boundary conditions on $\Omega = [0, 1]^2$
- Known smooth solution $\mathbf{U} \in H^2(\Omega)$
- Bilinear finite elements on uniform meshes, $h^{-1} = 8, 16, \dots, 128$
- Computations done using FOSPACK, by J. Ruge
- Lamé constants $\lambda = 2.15, \mu = 1.0$ (i.e., $\nu = 0.34$)



- Convergence of Newton's method
- Finite Element approximation rates
- Multigrid performance
- Nested Iteration efficiency

FE Discretization Accuracy

Nonlinear Functional: $\mathcal{G}(\mathbf{U}; f) = \|\nabla \cdot [(I + \mathbf{U})\Sigma(\mathbf{U})] - f\|^2 + \|\nabla \times \mathbf{U}\|^2$



- Optimal $O(h)$ convergence in H^1 and functional norms
- Optimal $O(h^2)$ convergence in L^2 norm

Computational Cost

Nested Iteration, 3 V(1,1)-pcg cycles per level

h^{-1}	m	$\mathcal{G}(\mathbf{U}^h; f)^{\frac{1}{2}}$	Ratio	$\bar{\rho}$	W_T	time (s)
8	1	2.64e-02		0.29	12.3	1
16	2	1.31e-02	2.01	0.25	16.0	4
32	3	6.61e-03	1.98	0.24	19.0	15
64	4	3.32e-03	1.99	0.24	20.8	60
128	5	1.66e-03	2.00	0.23	21.6	242

m Newton step

$\mathcal{G}^{\frac{1}{2}}$ Nonlinear functional norm

$\bar{\rho}$ Average V(1,1)-pcg convergence factor

W_T Total Cumulative Work: in # of Jacobi sweeps relative to current grid

Regularity in Nonsmooth Domains

Loss of smoothness at corners in polygonal domains is well understood (Kondrat'ev, Grisvard, Maz'ya, et. al.)

$$\mathbf{u} = \mathbf{u}_0 + \mathbf{S}(r, \theta)$$

$$\begin{aligned} \mathbf{u}_0 &\in H^2(\Omega) \\ \mathbf{S}(r, \theta) = r^\alpha \mathbf{\Theta}(\theta) &\in H^{\alpha+1}(\Omega) \end{aligned}$$

The value of α is determined by

- Boundary Condition Type: displacement, traction, mixed
- Geometry: interior angle at the corner
- Lamé Constants: λ, μ

This difficulty occurs in many second-order problems and is especially problematic for FOSLS discretizations!

Consider the FOSLS two-stage solution of Poisson's equation:

$$\left\{ \begin{array}{ll} \Delta p = f, & \Omega \\ p = 0, & \partial\Omega \end{array} \right. \implies \left\{ \begin{array}{ll} \nabla p - \mathbf{u} = 0, & \Omega \\ \nabla \cdot \mathbf{u} = f, & \Omega \\ \nabla \times \mathbf{u} = 0, & \Omega \\ p = 0, & \partial\Omega \\ \boldsymbol{\tau} \cdot \mathbf{u} = 0, & \partial\Omega \end{array} \right.$$

Consider the FOSLS two-stage solution of Poisson's equation:

$$\left\{ \begin{array}{ll} \nabla p - \mathbf{u} = 0, & \Omega \\ \nabla \cdot \mathbf{u} = f, & \Omega \\ \nabla \times \mathbf{u} = 0, & \Omega \\ p = 0, & \partial\Omega \\ \boldsymbol{\tau} \cdot \mathbf{u} = 0, & \partial\Omega \end{array} \right. \quad \begin{array}{l} G_1(\mathbf{u}; f) = \|\nabla \cdot \mathbf{u} - f\|^2 + \|\nabla \times \mathbf{u}\|^2 \\ G_2(p; \mathbf{u}) = \|\nabla p - \mathbf{u}\|^2 \\ \mathcal{V} = \{\mathbf{u} \in H^1(\Omega) : \boldsymbol{\tau} \cdot \mathbf{u} = 0 \text{ on } \partial\Omega\} \\ \mathcal{W} = \{p \in H^1(\Omega) : p = 0 \text{ on } \partial\Omega\} \end{array}$$

Consider the FOSLS two-stage solution of Poisson's equation:

$$\left\{ \begin{array}{ll} \nabla p - \mathbf{u} = 0, & \Omega \\ \nabla \cdot \mathbf{u} = f, & \Omega \\ \nabla \times \mathbf{u} = 0, & \Omega \\ p = 0, & \partial\Omega \\ \boldsymbol{\tau} \cdot \mathbf{u} = 0, & \partial\Omega \end{array} \right. \quad \begin{array}{l} G_1(\mathbf{u}; f) = \|\nabla \cdot \mathbf{u} - f\|^2 + \|\nabla \times \mathbf{u}\|^2 \\ G_2(p; \mathbf{u}) = \|\nabla p - \mathbf{u}\|^2 \\ \mathcal{V} = \{\mathbf{u} \in H^1(\Omega) : \boldsymbol{\tau} \cdot \mathbf{u} = 0 \text{ on } \partial\Omega\} \\ \mathcal{W} = \{p \in H^1(\Omega) : p = 0 \text{ on } \partial\Omega\} \end{array}$$

(1) Given $f \in L^2(\Omega)$, choose $\mathbf{u} = \underset{\mathbf{v} \in \mathcal{V}}{\operatorname{argmin}} G_1(\mathbf{v}; f)$

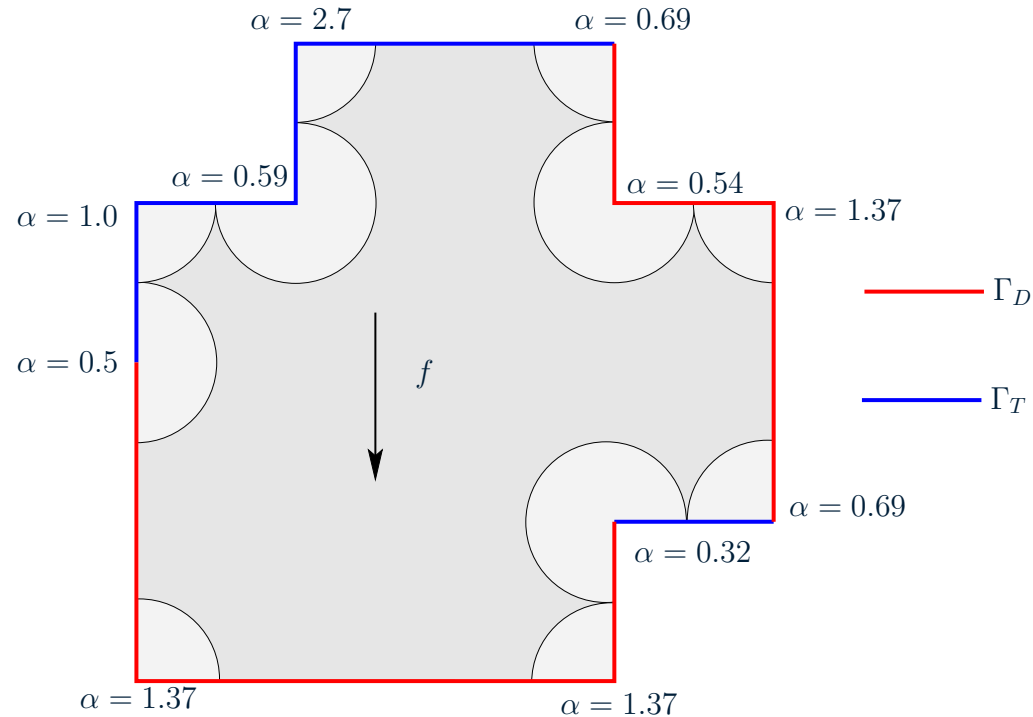
Consider the FOSLS two-stage solution of Poisson's equation:

$$\left\{ \begin{array}{ll} \nabla p - \mathbf{u} = 0, & \Omega \\ \nabla \cdot \mathbf{u} = f, & \Omega \\ \nabla \times \mathbf{u} = 0, & \Omega \\ p = 0, & \partial\Omega \\ \boldsymbol{\tau} \cdot \mathbf{u} = 0, & \partial\Omega \end{array} \right. \quad \begin{array}{l} G_1(\mathbf{u}; f) = \|\nabla \cdot \mathbf{u} - f\|^2 + \|\nabla \times \mathbf{u}\|^2 \\ G_2(p; \mathbf{u}) = \|\nabla p - \mathbf{u}\|^2 \\ \mathcal{V} = \{\mathbf{u} \in H^1(\Omega) : \boldsymbol{\tau} \cdot \mathbf{u} = 0 \text{ on } \partial\Omega\} \\ \mathcal{W} = \{p \in H^1(\Omega) : p = 0 \text{ on } \partial\Omega\} \end{array}$$

(1) Given $f \in L^2(\Omega)$, choose $\mathbf{u} = \operatorname{argmin}_{\mathbf{v} \in \mathcal{V}} G_1(\mathbf{v}; f)$

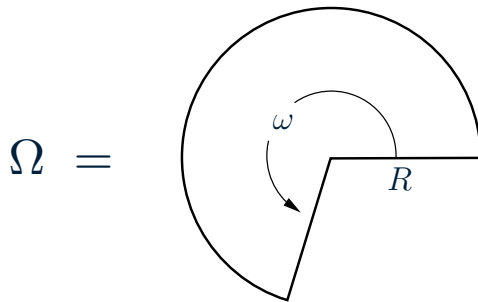
(2) Given $\mathbf{u} \in \mathcal{V}$, choose $p = \operatorname{argmin}_{q \in \mathcal{W}} G_2(q; \mathbf{u})$

Nonsmooth Domains: Example



- Each corner has a singularity of power α .
- Away from irregular boundary points, the solution is sufficiently smooth.

Nonsmooth Domains

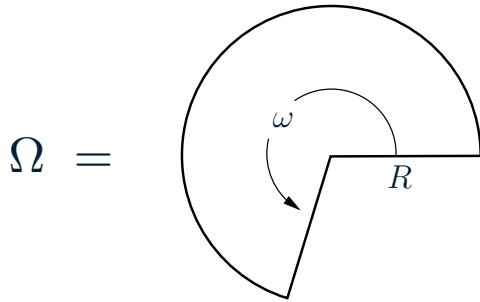


The solution to Poisson's equation near the corner may have a component of the form:

$$s(r, \theta) = r^\alpha \sin(\alpha\theta), \quad \alpha = \frac{\pi}{\omega}$$

(r, θ) a local polar coordinate system

Nonsmooth Domains



The solution to Poisson's equation near the corner may have a component of the form:

$$s(r, \theta) = r^\alpha \sin(\alpha\theta), \quad \alpha = \frac{\pi}{\omega}$$

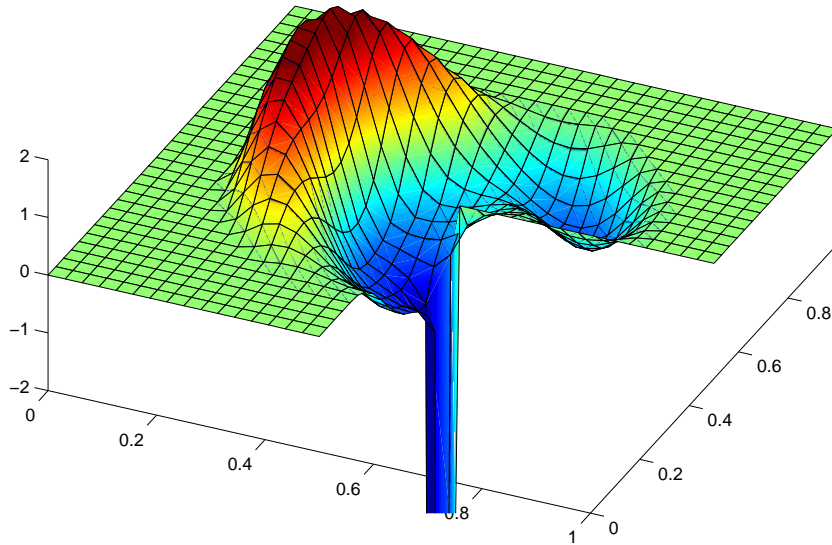
(r, θ) a local polar coordinate system

$$\begin{cases} p = p_0 + s & p_0 \in H^2(\Omega), \quad s \in H^{1+\alpha}(\Omega) \\ \mathbf{u} = \mathbf{u}_0 + \mathbf{S} & \mathbf{u}_0 \in H^1(\Omega), \quad \mathbf{S} \in H^\alpha(\Omega) \end{cases}$$

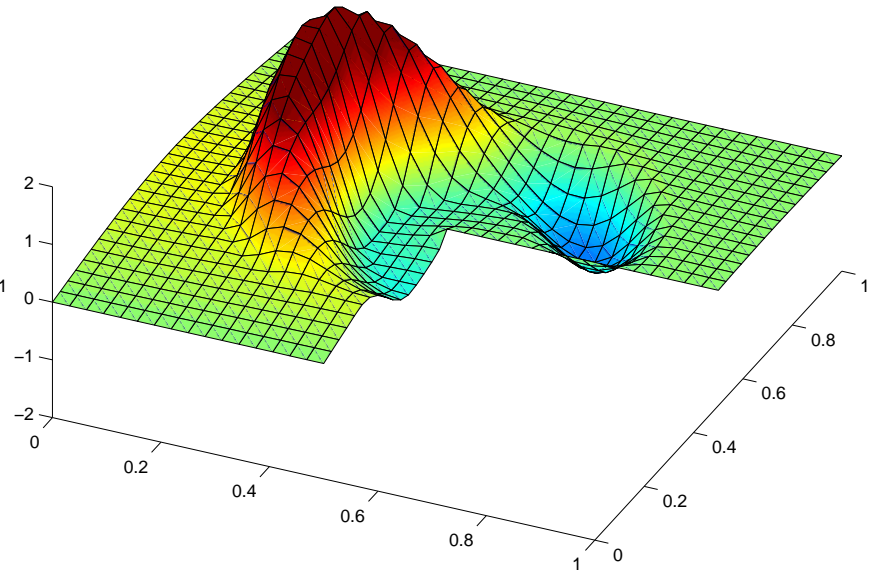
$$\omega > \pi \implies \alpha < 1 \implies \mathbf{u} \notin H^1(\Omega)$$

- Using standard FE subspaces of H^1 results in stalling convergence and a global loss of accuracy!

The Pollution Effect

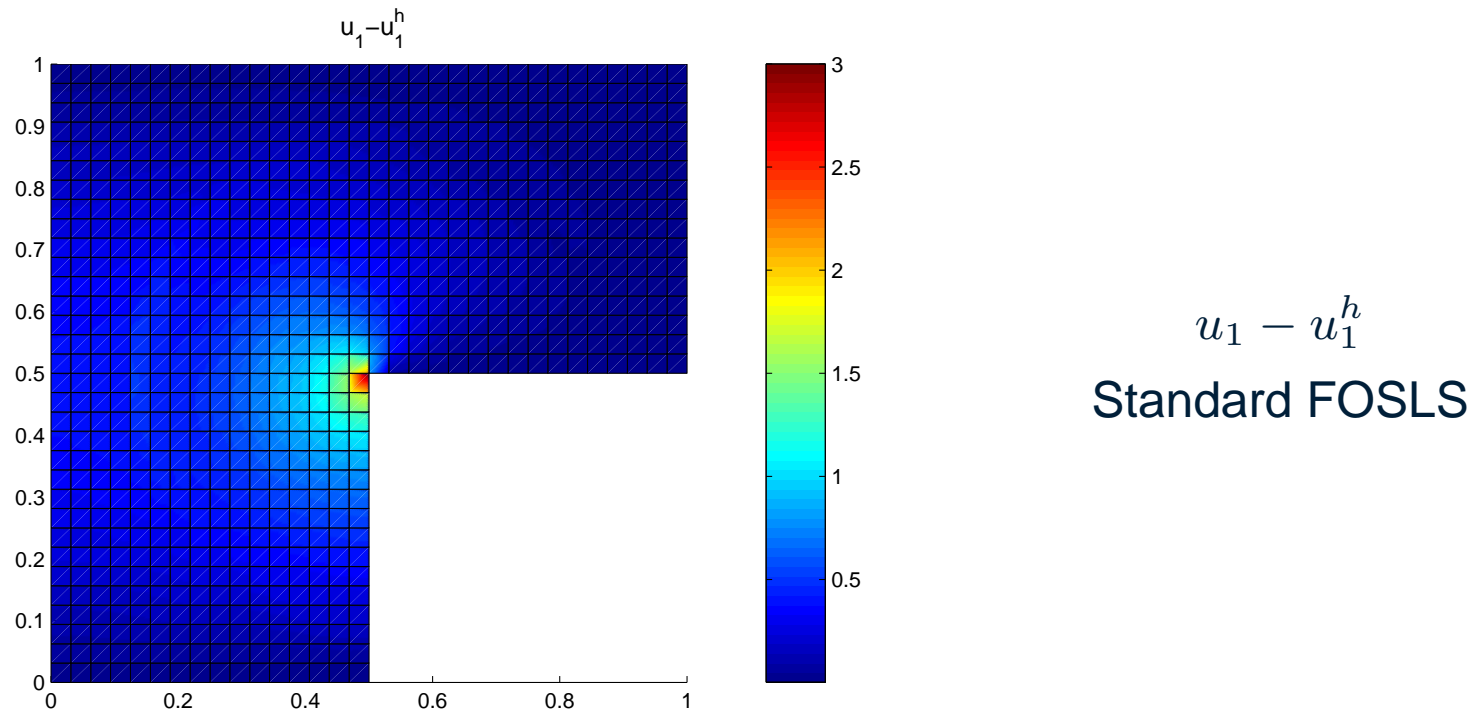


Exact component u_1



u_1^h computed by standard FOSLS

The Pollution Effect



- Even away from the singularity the solution is not accurate!

Existing Methods

- Add singular functions to finite element space
- Inverse-norm methods: H^{-1} least-squares, FOSLL*
- $H(\text{div})$ -conforming finite element spaces
- Weighted-norm FOSLS

Weighted-Norm FOSLS

Replace the L^2 norm functionals with locally weighted L^2 norms:

$$G_w(\mathbf{u}; f) = \|\mathbf{w}(\nabla \cdot \mathbf{u} - f)\|^2 + \|\mathbf{w}\nabla \times \mathbf{u}\|^2, \quad w = r^\beta$$

Let $\mathcal{V}^h \subset \mathcal{V}$ be the FE space. The discrete problem now becomes:

Given $f \in L^2(\Omega)$, choose

$$\mathbf{u}^h = \operatorname{argmin}_{\mathbf{v}^h \in \mathcal{V}^h} G_w(\mathbf{v}^h; f)$$

Choose the weight parameter β in order to

- make H^1 dense in $H(\operatorname{div}) \cap H(\operatorname{curl})$ in the weighted functional norm
- recover the best possible convergence throughout Ω

Define the weighted Sobolev norm and space:

$$\|\mathbf{u}\|_{k,\beta} = \left(\sum_{|\gamma| \leq k} \|r^{\beta+\gamma-k} D^\gamma \mathbf{u}\|^2 \right)^{\frac{1}{2}},$$

$$H_\beta^k(\Omega) = \{\mathbf{u} : \|\mathbf{u}\|_{k,\beta} < \infty\}.$$

for example:

$$\|\mathbf{u}\|_{1,\beta}^2 = \|r^\beta \nabla \mathbf{u}\|^2 + \|r^{\beta-1} \mathbf{u}\|^2$$

Define the weighted Sobolev norm and space:

$$\|\mathbf{u}\|_{k,\beta} = \left(\sum_{|\gamma| \leq k} \|r^{\beta+\gamma-k} D^\gamma \mathbf{u}\|^2 \right)^{\frac{1}{2}},$$

$$H_\beta^k(\Omega) = \{\mathbf{u} : \|\mathbf{u}\|_{k,\beta} < \infty\}.$$

Also define

$$L = \begin{pmatrix} \nabla \cdot \\ \nabla \times \end{pmatrix}, \quad \mathbf{f} = \begin{pmatrix} f \\ 0 \end{pmatrix}. \quad \implies \quad L\mathbf{u} = \mathbf{f} \quad \text{and}$$

$$G_w(\mathbf{u}^h; \mathbf{f}) = \|L\mathbf{u}^h - \mathbf{f}\|_{0,\beta}^2 = \|L(\mathbf{u}^h - \mathbf{u})\|_{0,\beta}^2$$

Theory

Let Ω be a polygonal domain with one reentrant corner of interior angle ω and assume $\mathbf{f} \in L^2(\Omega)$. Let $\mathbf{u} \in H^\alpha(\Omega)$ satisfy $L\mathbf{u} = \mathbf{f}$. If $\mathbf{u}^h \in \mathcal{V}^h$ is chosen to minimize the weighted functional,

$$G_w(\mathbf{u}^h; \mathbf{f}) = \min_{\mathbf{v}^h \in \mathcal{V}^h} \|L\mathbf{v}^h - \mathbf{f}\|_{0,\beta}^2,$$

for $1 - \alpha < \beta < \min(1 + \alpha, 2 - \alpha)$, then the approximation error, $\mathbf{u} - \mathbf{u}^h$, satisfies the following bounds:

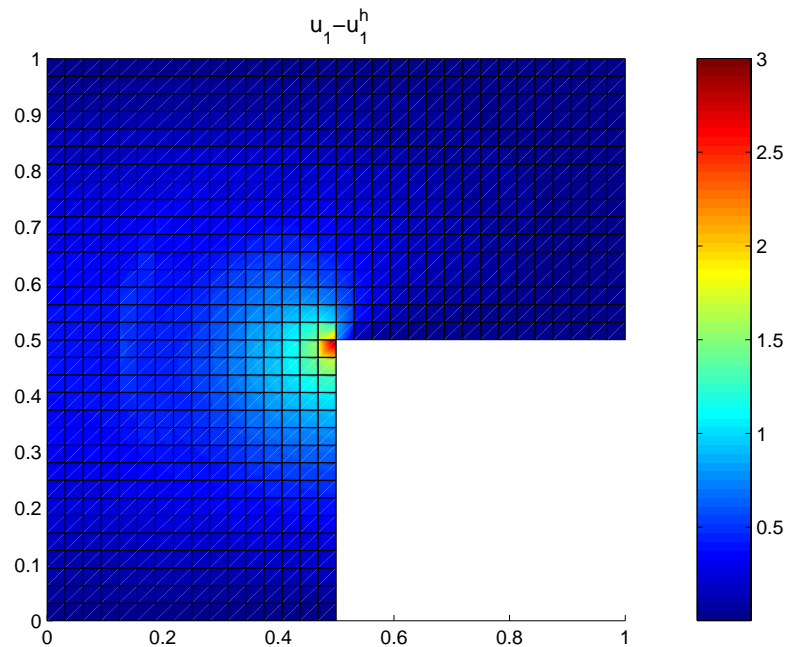
$$(1) \quad \|\mathbf{u} - \mathbf{u}^h\|_{1,\beta} \leq Ch^{\alpha+\beta-1} \|\mathbf{u}\|_{\alpha+\beta,\beta}, \quad O(h^{\alpha+\beta-1})$$

$$(2) \quad G_w(\mathbf{u} - \mathbf{u}^h; \mathbf{0})^{\frac{1}{2}} \leq Ch^{\alpha+\beta-1} \|\mathbf{u}\|_{\alpha+\beta,\beta}, \quad O(h^{\alpha+\beta-1})$$

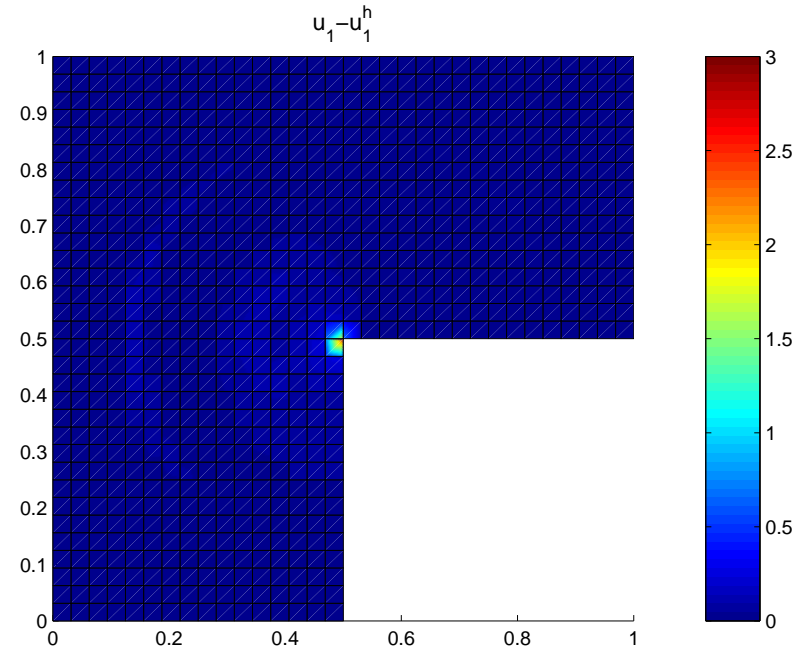
$$(3) \quad \|\mathbf{u} - \mathbf{u}^h\|_{0,\beta} \leq Ch^{\alpha+\beta} \|\mathbf{u}\|_{\alpha+\beta,\beta}, \quad O(h^{\alpha+\beta})$$

$$(4) \quad \|\mathbf{u} - \mathbf{u}^h\|_0 \leq Ch^\alpha (\|\mathbf{u}\|_\alpha + \|\mathbf{u}\|_{\alpha+\beta,\beta}). \quad O(h^\alpha)$$

Reduction of the Pollution Effect

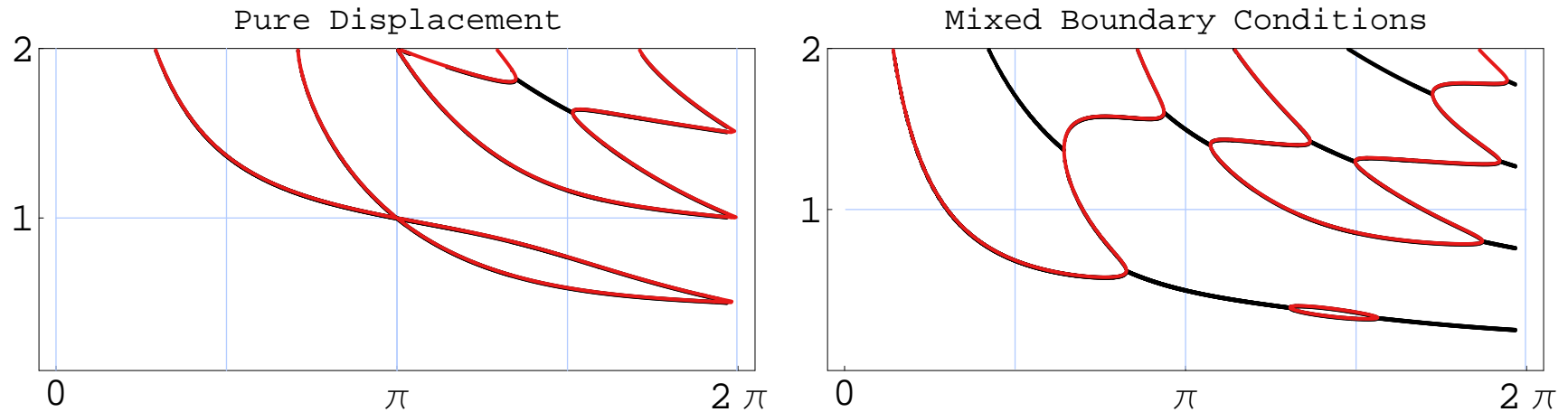


$u_1 - u_1^h$
Standard FOSLS, No weighting



$u_1 - u_1^h$
Weighted-norm FOSLS, $\beta = 4/3$

Powers of Singularities: Elasticity



Each plot is the power of the singularity α vs. interior angle ω for linear elasticity with Lamé constants $\lambda = 2.15$, $\mu = 1$.

- Pure Displacement $\omega \geq \pi \implies \mathbf{U} \notin H^1(\Omega)$
- Mixed Boundaries $\omega \gtrsim 0.3\pi \implies \mathbf{U} \notin H^1(\Omega)$

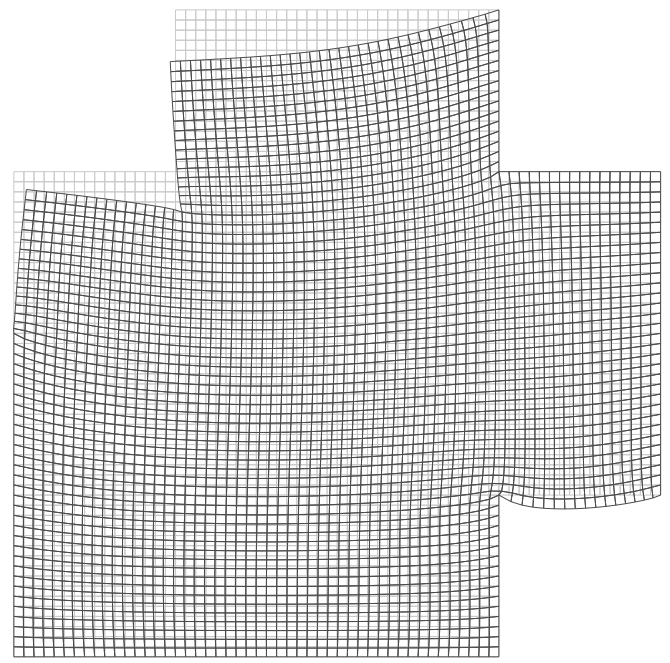
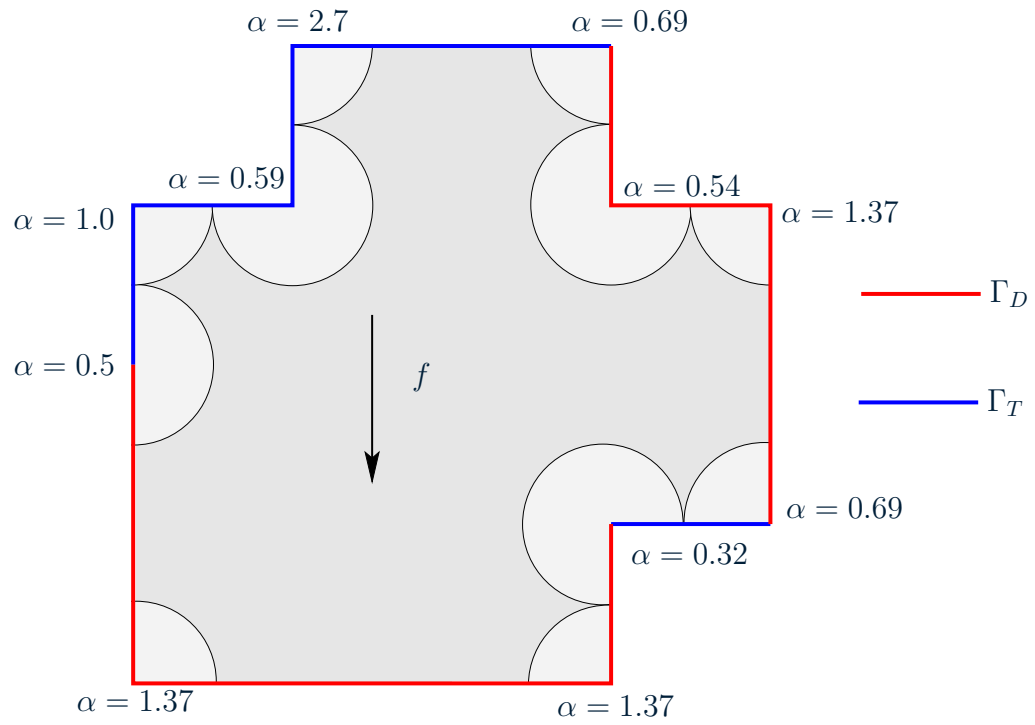
Numerical Results: Elasticity with Singular Solutions

- Mixed boundary conditions on Ω
- Interior body force $\implies \mathbf{U} \in L^2(\Omega) \setminus H^1(\Omega)$
- Use weighted functional:

$$G_w(\mathbf{U}; \mathbf{U}_n, \mathcal{F}_n) = \|\nabla \cdot \tilde{A}_n \mathbf{U} - \mathcal{F}_n\|_{0,\beta}^2 + \|\nabla \times \mathbf{U}\|_{0,\beta}^2$$

- Use weight of $\beta = 2 - \alpha$ near each corner

Domain and Boundary Conditions



$$\mathbf{U} \sim r^{\alpha-1} \in H^\alpha(\Omega)$$

Numerical Results: Singular Solutions

Nested Iteration, 5 V(1,1)-pcg cycles per level

h^{-1}	m	$\mathcal{G}_w(\mathbf{U}^h; f)^{\frac{1}{2}}$	Ratio	$\bar{\rho}$	W_T	time (s)
8	1	2.88e-01		0.49	20.1	7
16	2	1.94e-01	1.48	0.59	28.3	23
32	3	1.27e-01	1.53	0.76	33.0	90
64	4	8.11e-02	1.56	0.82	37.1	350
128	5	5.14e-02	1.58	0.82	38.9	1376

m Newton step

$\mathcal{G}_w^{\frac{1}{2}}$ Weighted nonlinear functional norm

$\bar{\rho}$ Average V(1,1)-pcg convergence factor

W_T Total Cumulative Work: in # of Jacobi sweeps relative to current grid

Concluding Remarks and Open Questions

Summary

- Robust solution strategy for geometrically-nonlinear elasticity problems in both smooth and nonsmooth domains.
- Corner singularities for many problems can be treated with weighted norms, restoring (nearly) optimal discretization convergence rates.

Open Questions

- Elasticity theory for pure traction, mixed boundary conditions
- Elasticity in the incompressible limit, $\lambda \rightarrow \infty$
- Improved multigrid solvers for weighted problems